

A project present to
The Faculty of the Department of Aerospace Engineering
San Jose State University

in partial fulfillment of the requirements for the degree

The Designated Masters' Project Committee Approves the Thesis Titled

* 'T+ +, AT+* - *F E))E - T.+) +T/ DU.+- 0 AT 1 * BU. - STAT+* - A) 2U+S+T+* -

SE2UE -)E *F A 0E*S / -) 3 . * - *US * .B+T

%y

4ynn)hou5

A ' ' . * 6ED F* . T3E AE . *S'A)E E - 0+- EE .+- 0 ' . * 0 . A

SA - J*S7 STATE U -+6E .S+T /

ay !"# \$

Dr. Teri (Lis) Papadopoulos

Aerospace Engineering

Dr. Sean S. Sei

- ASA Ames Research Center

Gary Johnson

- ASA Ames Research Center

&

Table #B . un AB Eccentricity 6arying First Burn Start Time!! !<

Table !B . un BB Eccentricity 6arying First Burn Start Time!! !=

' (

Figure #B Final Eccentricity 6arying Burn # and Burn !!!#&

Figure !B Final Drift 6arying Burn # and Burn !!!#<

Figure \$B *ptimal Burn Drift 'rofile!!#A

Figure <B *ptimal Burn Eccentricity 'rofile!!#=

Figure AB *ptimal Burn Drift 'rofile Targeting 4ongitude!!#>

Figure =B *ptimal Burn Eccentricity 'rofile Targeting 4ongitude!!#?

Figure >B *ptimal Burn 4ongitude 'lot!!#@

Figure ?B *ptimal Burn 4ongitude and Drift .ate 'lot!!!"

Figure @B . un AB Final Eccentricity 6arying Burn ! Start Times!!!#

Figure #"B . un AB Burn .atio 6arying Burn ! Start Times!!

Figure ##B . un AB Burn .atio 6arying Start Times!!!\$

Figure #!B . un AB Final Eccentricity 6arying Start Times!!!<

Figure #B \$B . un BB Burn .atio 6arying Start Times!!!A

Figure #<B . un BB Final Eccentricity 6arying Start Times!!! =

" &

A Unperturbed geostationary semimajor axis $C = \frac{4\pi^2 a^3}{GM}$ (m)

a * r%it semimajor axis

D Drift rate

e * r%it eccentricity

E)F Earth entered Fixed

E)+ Earth entered Inertial

0AST 0reen5 ich Apparent Sidereal Time

0 ST 0reen5 ich Mean Sidereal Time

i inclination

p Semilatus rectum

'osition vector of the satellite

r magnitude of the position vector of the satellite

\dot{r} Acceleration vector

raan . ight ascension of the ascending node

) 6elocity vector

6 magnitude of velocity vector

E Argument of perigee

v True anomaly

; ; component of position vector

y y component of position vector

8 8 component of position vector

* +

Optimal maneuvers have been studied for a long time. Hohmann solved the problem of transfers between two coplanar circular orbits with a minimum velocity applied to the space vehicle. Hohmann's work! The problem of two impulse orbit maneuvers has been studied over the years with minor breaks in order to apply solutions to real spacecraft. Jespersen and Littleman's work! wrote about an analytical approach to two fixed impulse transfers. Shilev and Melton's work! concluded that using two impulsive maneuvers of fixed magnitudes is only possible for certain thrust directions.

In the current age of affordability in the Aerospace industry, optimization problems are becoming increasingly important. As a satellite nears completion of orbit transfer, plans must be made in order to place the satellite into its final orbit. Geostationary satellites are equipped with thrusters that allow the satellite to be commanded to maneuver the spacecraft into the desired orbit. These burns are usually tangential to the orbit plane or orthogonal to the orbit plane. Tangential burns are also known as along-track burns, and change the longitude of the satellite. This affects the semimajor axis, the longitude drift rate and the eccentricity vector. Orthogonal burns change the orientation of the orbit plane. This includes the inclination and the ascending node. Soop's work!

This paper deals only with tangential burns in line with the orbit plane. In order to stop at the final station location in geosynchronous orbit, the satellite needs to have a semimajor axis of 42164 km (a drift rate of zero) and an eccentricity vector very close to zero.

It is usually optimal to command orbit maneuvers during the apsides of the orbit. Sguini, M. Teofilatto's work! put with affordability being a main driver, for the sake of time, orbit maneuvers are sometimes scheduled at non-optimal times. The time of the maneuver can be

varied to change the eccentricity vector direction, however, with this being set by limitations of time, having multiple burns enables some variability in the final eccentricity of the orbit. This paper will deal with two burns separated by 1 hour. The first objective of this work is to identify the sizes of these two burns for a specific orbit that will give the result of a final eccentricity as close as possible to the target eccentricity. The second objective is to find any general conclusions between the burn sizes and eccentricity that can be used on general orbits.

, % &

The investigation was done by varying the ratio of the turns and propagating through an ephemeris at a specified turn time using Matlab as the computing tool. The orbit was modeled using the two-body model equations with no perturbations. The input sheet includes the ephemeris in classical orbit parameters which are then rotated into Earth centered inertial (ECI) coordinates for propagation using the equations of Mars.

Equation 1

$$p = a(1 - e^2)$$

Equation 2

$$r = \frac{p}{(1 + e \cos \theta)}$$

Equation 4

$$\begin{aligned}
 & + \dot{\theta} \\
 & + \dot{\theta} \\
 & (\dot{\theta} + e \cos \theta) \\
 & \cos \theta \\
 & \dot{\theta} \\
 & + \dot{\theta} \\
 & + \dot{\theta} \\
 & (\dot{\theta} + e \cos \theta) \\
 & \cos \theta \\
 & \dot{\theta} \\
 & + \dot{\theta} \\
 & \cos(\dot{\theta} + e \cos \theta) \\
 & i(\dot{\theta}) \\
 & (\dot{\theta} + e \sin \theta) - \cos \theta \dot{\theta} \\
 & \sin \theta \\
 & \sin \theta \dot{\theta} \\
 & (\dot{\theta} + e \sin \theta) + \sin \theta \dot{\theta} \\
 & \sin \theta \\
 & \cos \theta \dot{\theta} \\
 & \dot{\theta} \\
 & - \frac{\mu}{p} \dot{\theta} \\
 & V = \begin{bmatrix} x \\ y \\ z \end{bmatrix} = \dot{\theta}
 \end{aligned}$$

The vector can be propagated forward in time by multiplying a time step with the velocity and acceleration found from equation A16.

Equation 5

$$\dot{r} + \frac{\mu}{a^3} r = 0$$

The initial drift rate of the orbit or how quickly the orbit is rotating with respect to the rotation of the earth is given by equation A17. To find the thrust needed to reduce the drift rate to zero, Equation A18 provides equation 5.

Equation 6

$$D = \frac{-1.5}{A} a$$

#

Equation 7

$$D = \frac{-3 \text{ V}}{V}$$

Using the F6, an array is created by splitting the value between the two turns by a percentage. In this simulation, one percent is used as the step change between each case. A custom Runge-Kutta algorithm is written in order to propagate the ephemeris while controlling the time steps. In the simulation a time step of 10 seconds is used. The Runge-Kutta algorithm also checked to see if the turn time is passed during the calculation step, and if it is, inserted an instantaneous thrust or delta V at that time.

In order to see if the space vehicle ends up at the correct station, longitude also needed to be calculated. Although the ephemeris is already calculated in Earth-centered Earth-fixed (ECEF) coordinates, this needs to be converted to an Earth-centered inertial (ECI) coordinate system. Since longitude is only based on the position, we only need to convert the position vector from ECEF to ECI given by Eagle and J in equation 7 and 8.

Equation 8

$$r_{ecf} = [T] r_{eci}$$

Equation 9

$$[T] = \begin{bmatrix} \cos \theta & \sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

where θ is equal to the Greenwich sidereal time at the moment. This can be calculated from the current day and time by first calculating the Julian Date using Equation 16 (Allado et al., 2009), then by finding Greenwich Mean Sidereal Time (GMST) by using equation 17 and 18 (Allado et al., 2009).

Equation 10

$$\text{JD} = 367 \text{ year} - \left\{ \frac{7 \left\{ \text{year} + \left(\frac{\text{month} + 9}{12} \right) \right\}}{4} \right\} + \left(\frac{275 \text{ month}}{9} \right) + \text{day} + 1721013.5 + \left(\frac{\text{second}}{60} + \text{minute} \right)$$

The equations for nutation in longitude and nutation in obliquity using the trigonometric arguments are listed below in arc seconds (Eagle and J)

Equation 18

$$\Delta\lambda = -17.20 \sin \lambda' - 1.32 \sin 2\lambda' - .23 \sin 2\lambda'' + 0.21 \sin 2\lambda'''$$

Equation 19

$$\Delta\epsilon = 9.20 \cos \lambda' + 0.57 \cos 2\lambda' + 0.10 \cos 2\lambda'' - 0.09 \cos 2\lambda'''$$

After converting the results from equations 18 and 19 into degrees and putting them along with the result from equation 17 into equation 15 the Greenwich Apparent Sidereal Time in degrees is obtained. This can be put into equation 14 which will create the matrix; to convert the position vector into the position vector.

Equations 18 and 19 provide the equation to calculate longitude from the position vector.

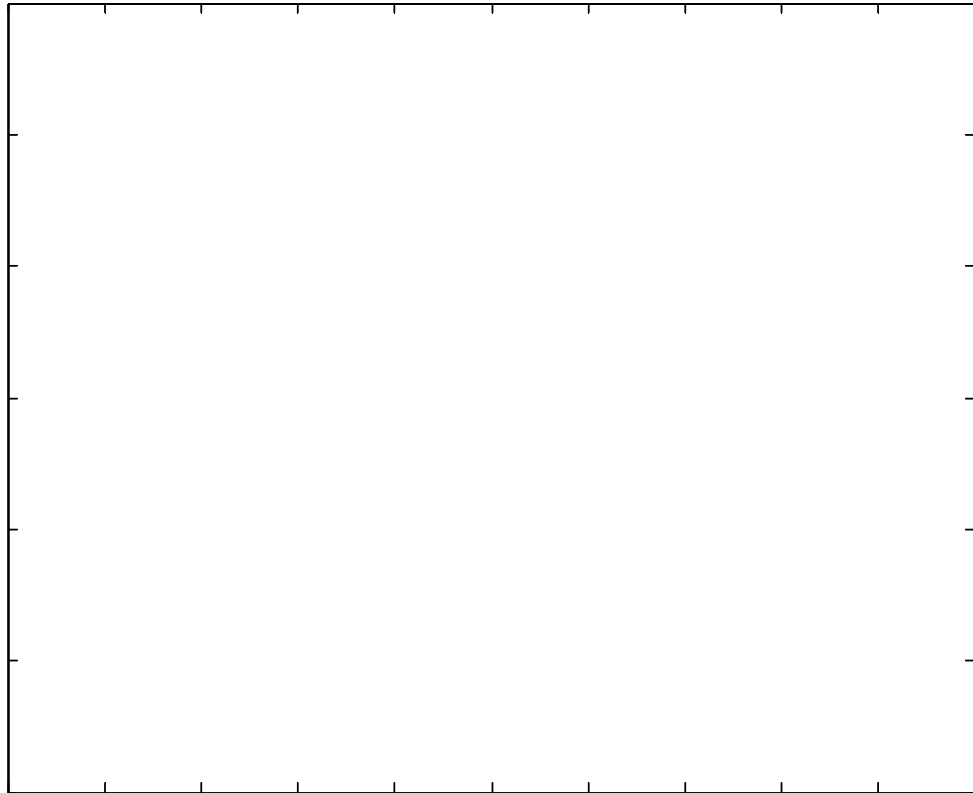
Equation 20

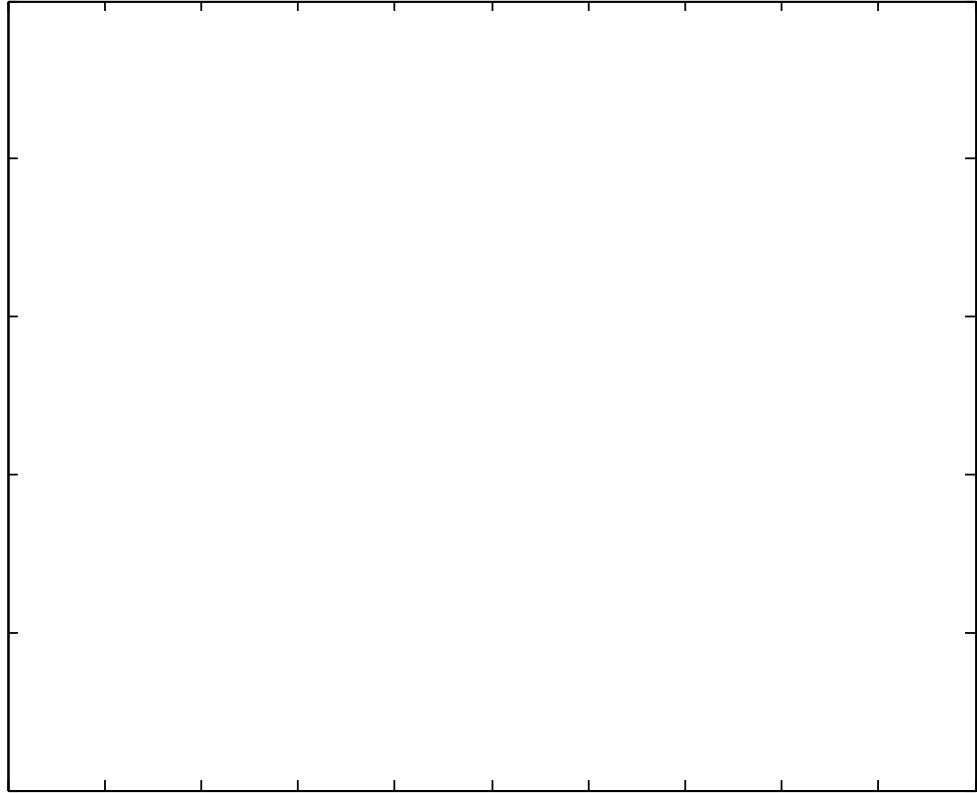
$$\lambda = \tan^{-1} \frac{y}{x}$$

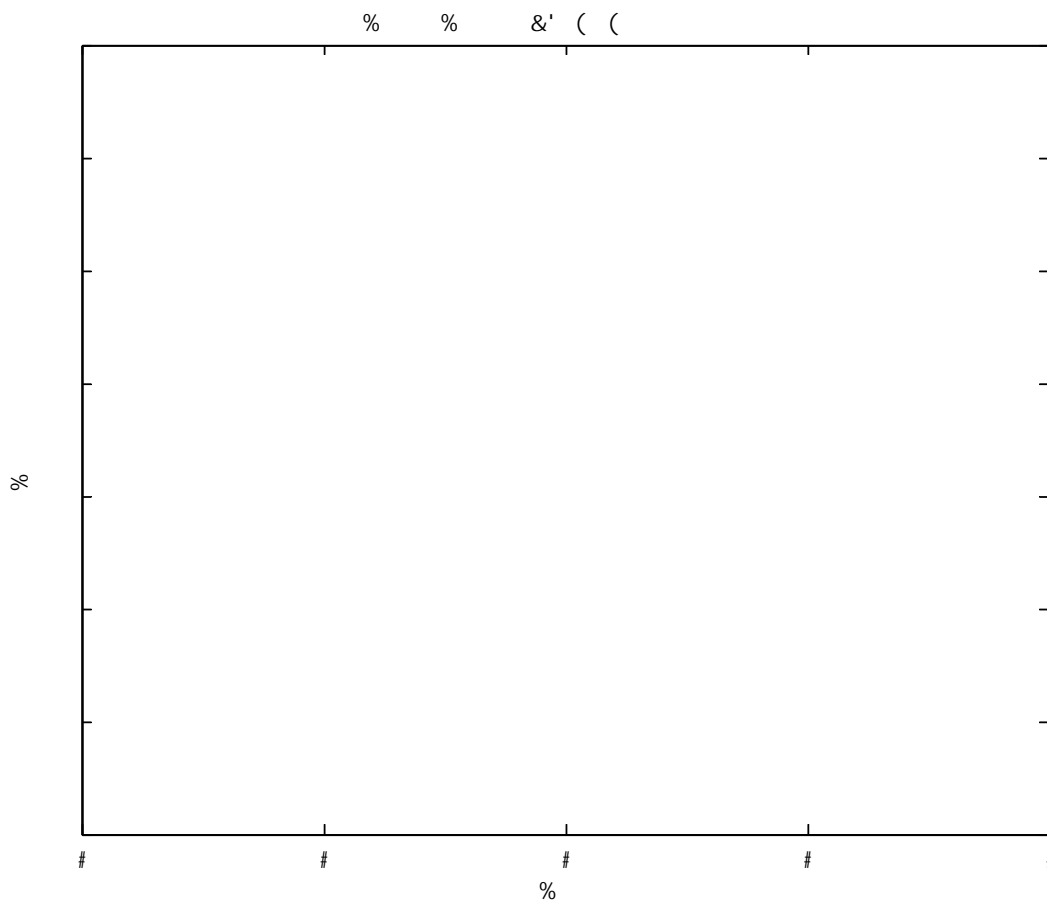
The longitude value is calculated at every step of the range Nutta and at the end of the script the parameter with the longitude closest to the desired longitude is selected.

- &

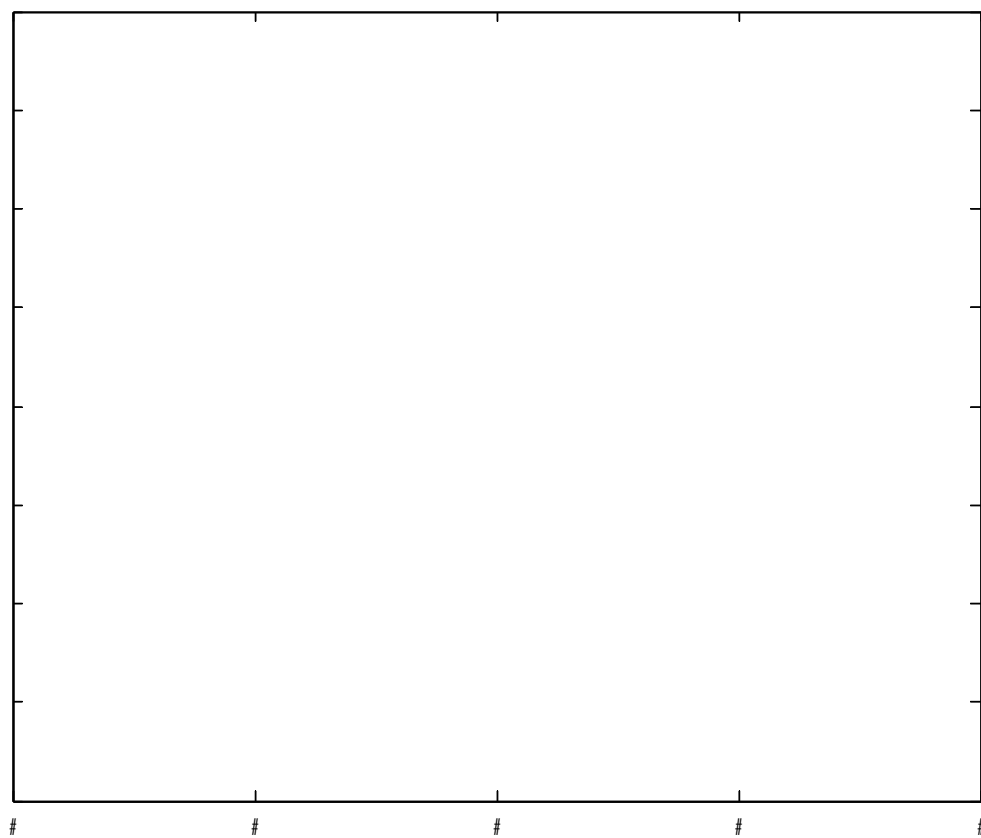
For this simulation, an orbit and a burn time were randomly selected. A starting drift rate of # deg/day least 5 and J 5 as selected as this seems to be reasonable for current satellites. The input file can be seen in Appendix A. These results do not target a specific longitude.



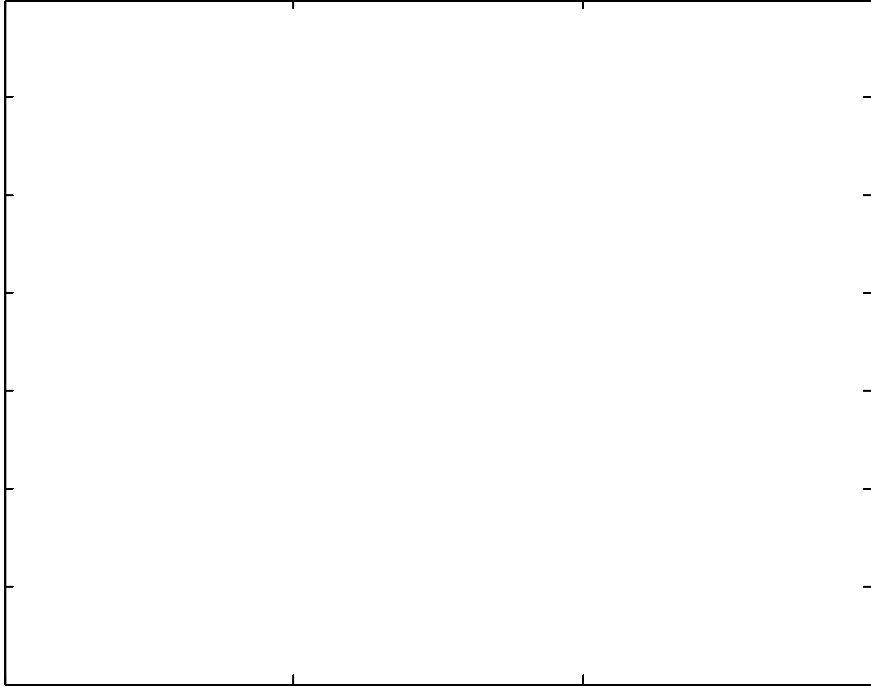




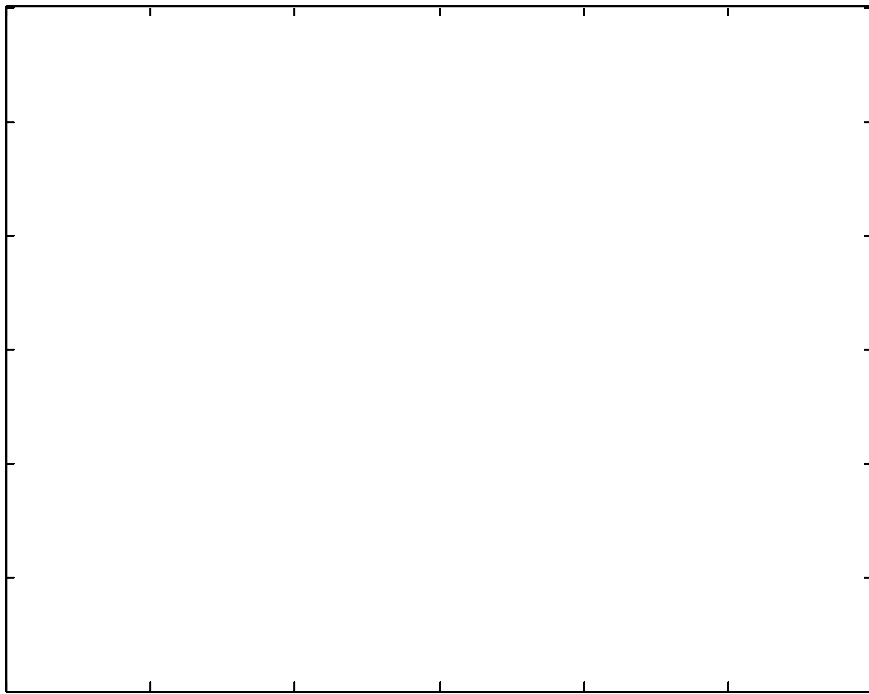
'igu e 3(, E713abPou 4) 57a9\$ n1t5B7u9 5%#2qfa9b7D\$

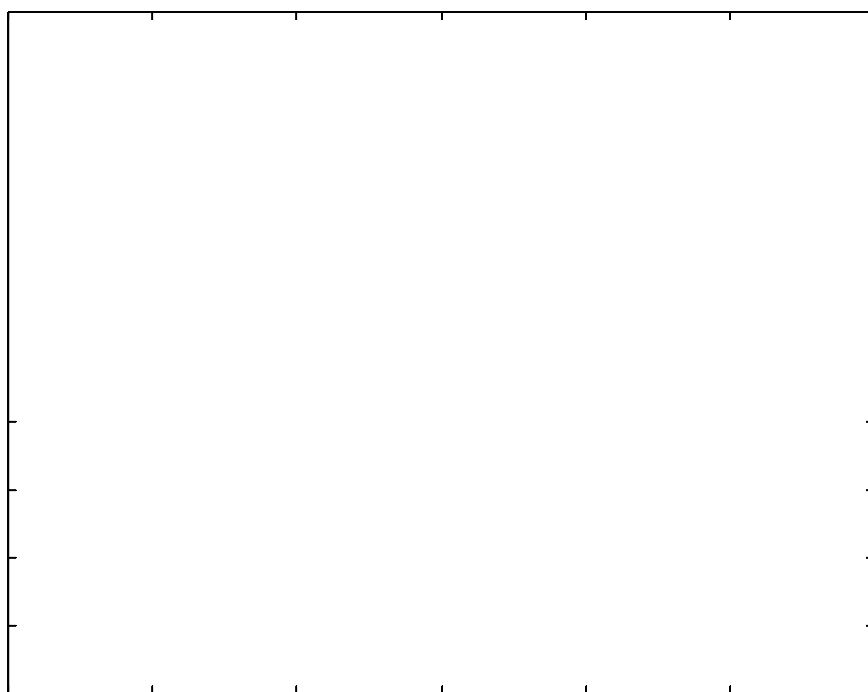


In comparison to Figure 9, Figure A shows an optimal ratio of λ . This is very different from the ratio of λ for Figure 9, however, both still result in a drift rate close to zero.



second $\%urn$ and its result for $\%urn$ ratio and eccentricity $\&$





Run A: Ratio vs Burn 1 Start Time

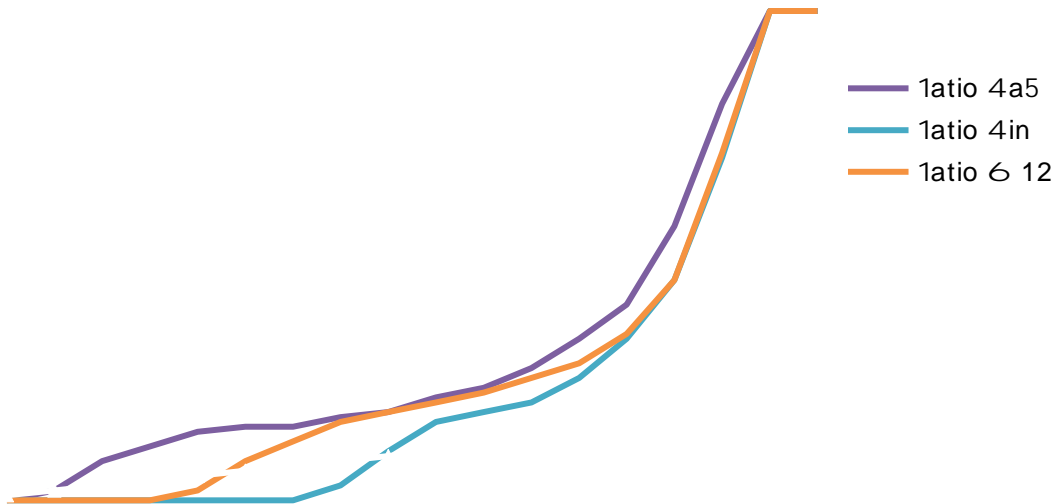


Figure 11 (Burn 2 (Burn 1 ratio) starting times) plots the optimal ratio for the first burn shown by the horizontal axis and the second burn # hours later. The ratio in and a; show the minimum and maximum optimal ratios when varying the start time of the second burn by three hours.

There is a limit on both extremes. The values all go to zero when the burns start too early. Since the simulation calculates an optimal d_6 based on one impulse, if the burn starts when the initial longitude is too far away there is no way that the d_6 will be able to achieve the desired longitude with any burn ratio. Likewise if the first burn starts too late, there is no time for a second burn and the burn ratio is always # " " R.

Run A: Ecc vs Burn 1 Start Time

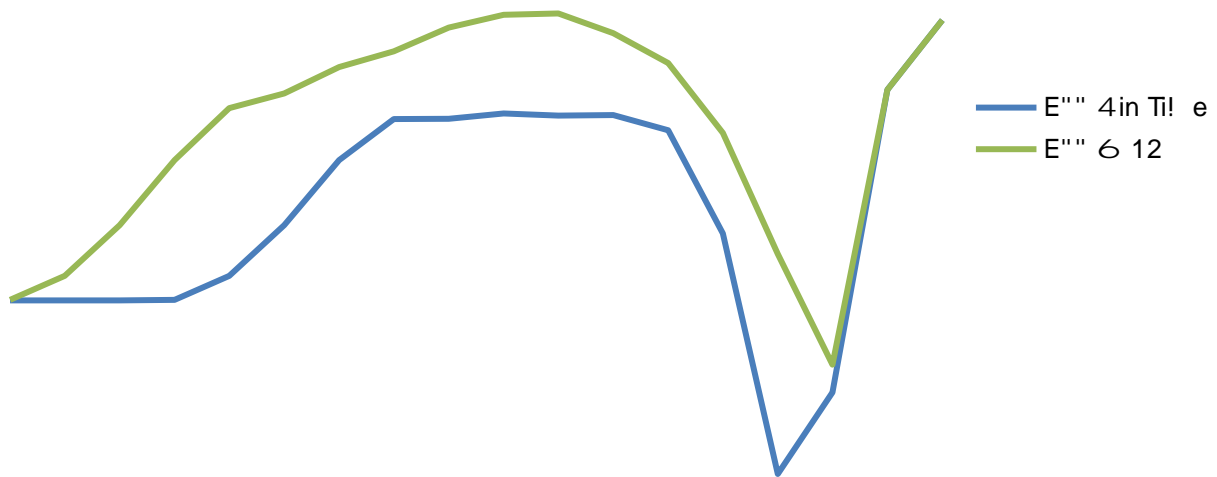
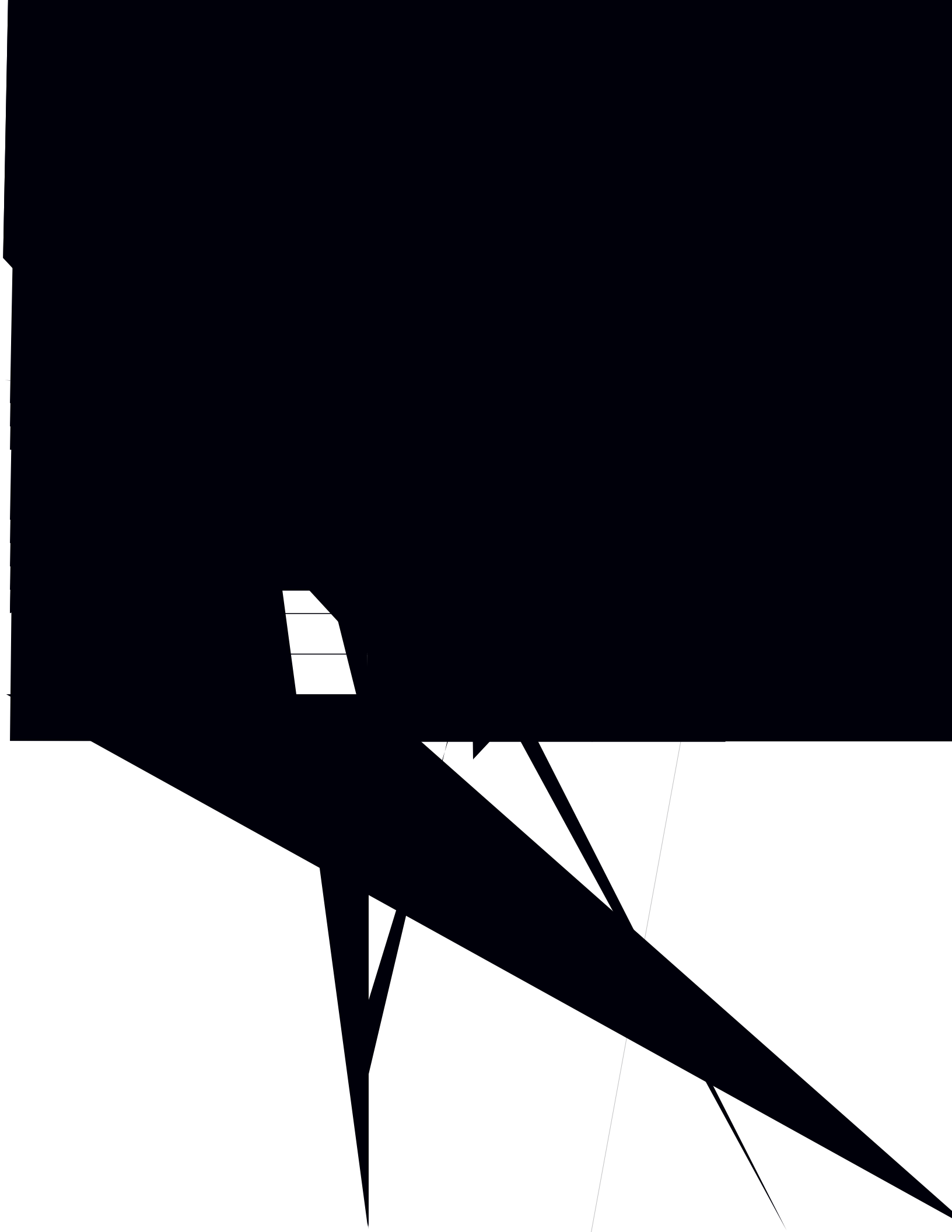
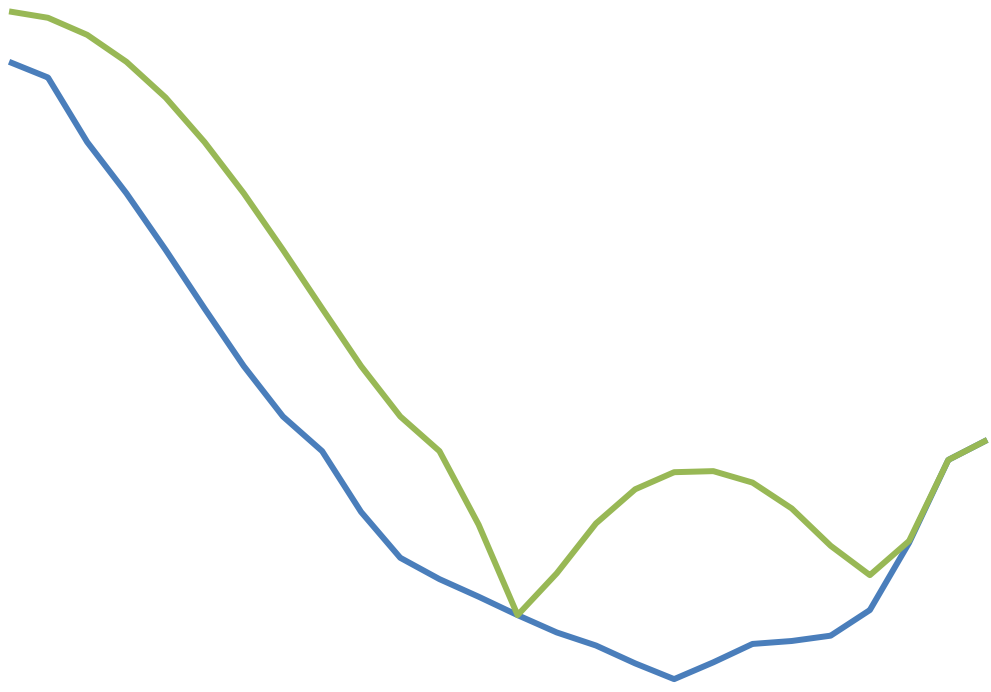


Figure 12(1) shows the final eccentricity (e) as a function of the start time of the second burn. Similar to the Burn Ratio plot shown in Figure 9, Ecc S #! plots the final eccentricity for the first burn shown by the horizontal axis and the second burn #! hours later. Ecc in Time shows the minimum final eccentricity when varying the start time of the second burn by three hours.

Ecc S #! is consistently larger than the minimum eccentricity e;cept near the end of the data. Table # shows the reaction of the time of the minimum eccentricity.







| | | | |
|---------------|---------------|-------|-----------|
| | 75 | | |
| 4728713 11(00 | 0%00062 67 | 11%25 | 0%0006334 |
| 4728713 12(00 | 0%00095 8 | 12 | 0%000958 |
| 4728713 13(00 | 0%00103 8 | 12 | 0%001038 |

4i(e in . un A) the eccentricity is at a minimum near 5 hen the %urn ratio is " R and the time %et5 een the %urns is high. The optimum time %et5 een %urns also decreases as the %urn ratio gets higher.

• &

0iven an initial or%it and %urn time, this simulation successfully ans5ers the question of 5 hat si8e t5 o %urns #! hours apart should %e in order to get a minimum eccentricity value 5 hile targeting a specific longitude. In addition, the simulation can vary the separation time %et5 een the t5 o %urns in order to further minimi8e the eccentricity value.

1 hile this simulation does ans5er the specific question of the optimal %urn percentage split for a particular or%it, it is very specific to each case. A general conclusion that can %e made is that if the d6 needed to get 8ero degrees drift rate is too small, no amount of varying the %urns

satellite to point in the correct direction. Also, as satellites reach their final geosynchronous orbit, they start to keep track of time in terms of the satellite's position and the sun's position. This allows spacecraft operators to keep track of when the satellite will be subject to thermal constraints. Adding this time system to the simulation will make it easier for users to calculate burn times.

0

. un A input file

| | |
|-------------------------|---------------------|
| +nitial Semi major A;is | B <!?"?=" |
| +nitial eccentricity | B &" "" "\$#<<<\$A? |
| +nitial inclination | B &##!A<@=@! |
| +nitial raan | B D&=@#!?<=A> |
| +nitial arg of perigee | B \$&#\$?A>##"## |
| +nitial true anomaly | B D" &#?!<<!>A> |
| Solar . adiation I) pJ | B #&!A |
| +nitial epoch year | B !"#! |
| +nitial epoch month | B "\$ |
| +nitial epoch day | B !? |
| +nitial Epoch hour | B #! |
| +nitial Epoch minutes | B # \$ |
| +nitial epoch seconds | B \$ # |
| Burn epoch year | B !"#! |
| Burn epoch month | B "\$ |
| Burn epoch day | B !@ |
| Burn Epoch hour | B "" |
| Burn Epoch minutes | B "" |
| Burn epoch seconds | B "" |

0

. un A input file

+nitial Semi major A; is I(mJ B <! " ? =
+nitial eccentricity B & " " " @

